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## Highly-crystalline $\gamma$ -MnS nanosaws†

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We report here a straightforward method to fabricate single-crystal wurtzite-structured manganese sulfide nanosaws ( $\gamma$ -MnS NSs) via a chemical vapor deposition technique, which can be extended to different types of binary chalcogenide-based 1D building blocks. Advanced spectroscopy and electron microscopy techniques were applied to reveal the structure, morphology and composition of these nanostructures. Our studies indicate that  $\gamma$ -MnS NSs are of high crystalline quality and purity, exhibit high aspect ratio and have extrusive teeth with average tip size of approximately a couple of atoms. These NSs are structurally stable and do not present divergence in surface energy. The formation of the backsaw is governed by a vapor–liquid–solid growth mechanism, whereas the teeth are a result of a self-catalyzed growth process induced by the Mn-terminated (0001) surface. Our optical and electrochemical analyses indicate that  $\gamma$ -MnS NSs can be used as efficient visible light emitters with an activation energy of  $\sim 40$  meV, and as anode materials with improved cycling stability for LIBs. This bottom-up approach represents a step ahead to provide promising and potential Mn-based materials for high-performance LIB technology and optoelectronics, and adds a new member to the family of saw-like wurtzite-structured nanoribbons, such as CdSe and ZnS.

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## Introduction

One-dimensional (1D) new-generation building blocks, such as nanowires (NWs), nanobelts (NBs), nanoribbons (NRs) and nanotubes (NTs), are appealing for understanding the basic science behind dimensionality, morphology-tuned stability, quantum-confinement transport phenomena and transformation mechanisms, and for device elements and interconnects.<sup>1–9</sup> Binary semiconductor NWs and NTs, which exhibit cylindrical symmetric cross sections, present a higher degree of functionality when compared to element semiconductors.<sup>4</sup> In the last past decade, single-crystal NBs of semiconducting oxides, possessing rectangular cross-sections and high surface area, have been employed for field-effect transistors, gas sensors, lithium-ion batteries, nanocantilevers and nanoresonators, and they have been found to exhibit an excellent performance.<sup>10–12</sup> NBs of semiconducting sulfides, such as ZnS, have been fabricated by different approaches for field emitters, ultraviolet-light sensors and photoconductive semiconductor switches.<sup>13–15</sup> In addition, it was also reported that ZnS NBs with specific morphology (wurtzite hexagonal structure) exhibit

enhanced structural stability, *i.e.*, they possess a more favorable low-energy surface structure.<sup>16</sup> Nonetheless, the paucity of wurtzite-type sulfur-based NBs limits the extended development of technological applications.

Manganese sulfide (MnS) is a p-type semiconductor ( $E_g \sim 3.2$  eV at room temperature) and crystallizes in three distinct phases: rock-salt ( $\alpha$ -MnS), zinc-blende ( $\beta$ -MnS) and wurtzite ( $\gamma$ -MnS).<sup>6</sup> Metastable hexagonal bulk  $\gamma$ -MnS does not exist in nature, but its production in the laboratory under unusual conditions shows that  $\gamma$ -MnS exhibits unique chemical, electrical and magneto-optical properties superior to those of stable  $\alpha$ -MnS.<sup>17,18</sup> Currently, studies on the architectural control of nanostructured  $\gamma$ -MnS have revealed that these nanostructures present high photoluminescence, improved antiferromagnetic response, and high reversible capacity and cycling stability.<sup>6,19–23</sup> Such advanced characteristics are expected to be exploited in improving the performance of optoelectronic devices, in developing new elements for diluted magnetic semiconductor (DMS) technology, and in implementing a novel anode material for Li-ion batteries (LIBs). Even though some reports on the synthesis of  $\gamma$ -MnS nanofibers have appeared in the literature,<sup>24–26</sup> studies focused on the fabrication of  $\gamma$ -MnS have failed to produce single-crystal nanosaws, *i.e.*, nanoribbons exhibiting one side flat and the other side saw-toothed. Their single-crystal nature are critical to avoid the grain boundary scattering present in polycrystalline MnS that affects the optoelectronic performance due to the high density of nanosized grains distributed across the surface. Also, their striking morphology would be crucial for any desired, potential and promising application.

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In order to provide a new 1D building block based on MnS having a non-conventional cross-section, a higher surface-to-volume ratio (when compared to nanorods of MnS), and a highly-stable surface structure, we report here for the first time the fabrication of single-crystal wurtzite-type MnS nanosaws (NSs) *via* a chemical vapor deposition (CVD) route that is both inexpensive and environmentally-friendly. This particular morphology opens up the possibility of extending the scope of applications of  $\gamma$ -MnS, including: customized nanocantilevers for probing scanning technology, gas sensors, field-effect transistors, nanoresonators, photoconductive semiconductor switches, visible light emitters and anode materials for high-performance LIBs. In this regard, we also investigate the photoluminescence (PL) properties of  $\gamma$ -MnS NSs and evaluate their performance as an anode material for LIBs.

## Experimental techniques

### Materials

All reagents used in this investigation, manganese chloride ( $\text{MnCl}_2$ ), gold chloride ( $\text{AuCl}_3$ ), and sulfur powder, were analytical-grade reagents (>99.999%) from Sigma Aldrich, USA, and were directly used without any further processing.

### Fabrication of $\gamma$ -MnS nanosaws

$\gamma$ -MnS NSs were grown on (001)Si substrates, which were thoroughly cleaned with acetone and then dried with nitrogen.  $\text{AuCl}_3$  was dissolved in isopropanol, and then ultrasonicated for 30 min. This homogeneous solution was next dripped onto the cleaned Si-substrate surfaces achieving a full and uniform coverage, and heated up to evaporate the solvent and any remaining absorbate. The Au catalyst-covered Si substrates,  $\text{MnCl}_2$  and sulfur powder were placed into a fused silica tube in a hot-wall horizontal furnace. The amounts of sulfur and  $\text{MnCl}_2$  were adjusted in order to obtain a 1 : 3 Mn/S molar ratio.  $\text{MnCl}_2$  was positioned in the middle of the furnace (hot zone), where the temperature ( $1273 \pm 10$  K) was monitored; while the sulfur (downstream) and the Si substrate (upstream) were placed at certain distances from the  $\text{MnCl}_2$  (170 and 115–100 mm, respectively) and kept at  $\sim 673$  K and 973–1123 K, respectively (see (ref. 6) for a depiction of the apparatus employed). The reaction tube was first evacuated for several hours down to 465 mPa to purge oxygen, heated to 1273 K, and filled to 46.5 Pa with a constant Ar flow at 20 sccm, which was used as carrying gas to transport the sublimated vapor to the cooler regions within the chamber for deposition. It was observed that after 1 h, the  $\text{MnCl}_2$  and S were fully consumed. After this, the tube was left to cool spontaneously to room temperature. In the absence of catalytic material, no  $\gamma$ -MnS NSs were obtained. The nanosaws were transferred to TEM grids upon carefully scratching the Si surface using a diamond tip for characterization and analyses.

### Characterization

The crystalline structure of the samples was studied using an X-ray diffractometer (XRD), Model Siemens D5000 with  $\text{Cu K}\alpha$  radiation. Raman spectra were collected *via* a Jobin-Yvon

T64000 spectrometer with Ar-ion laser excitation (514.5 nm) and attached to an optical microscope with  $80\times$  resolution. The surface morphology, elemental composition and topography were analyzed using a JEOL JSM-7500F field-emission scanning electron microscope (FE-SEM) and a JEOL JEM-2200FS Cs-corrected high-resolution TEM (HRTEM). All the images were taken in scanning TEM (STEM) mode. To collect temperature-dependent PL spectra, the samples were first mounted on the cold-finger of a closed-cycle helium-cryostat operating from 10 to 300 K under high vacuum condition ( $<0.1$  mPa). The samples were then excited with a 325 nm He–Cd laser (Model IK 3202R-D) with a maximum power of  $10 \text{ mW mm}^{-2}$ . The PL signal of the samples was dispersed by a 0.3 m Acton spectrometer with changeable holographic gratings and analyzed by a Princeton Instrument PI-MAX CCD camera equipped with UV intensifier, operating in the spectral region 200–950 nm. The decay curves were obtained by exciting the samples with an intensity-modulated He–Cd laser beam using an acousto-optic light modulator (Model AOM-125U) with a maximum modulation frequency of 125 MHz. The PL decay signals were detected by a photomultiplier (Hamamatsu R316-02) and a photon counting system with a turbo-multichannel scaler (Turbo-MCS, EG&G, Ortec), which offers a wide range of channel dwell time (minimum 5 ns) with no dead time between channels. The photomultiplier was equipped on an HR320 spectrograph (0.32 m) with a Czerny Turner configuration. The electrochemical characterization has been carried out in coin cells with Li as the counter and reference electrode, and  $\gamma$ -MnS NSs as the working electrode material. The anodes were fabricated by mixing  $\gamma$ -MnS nanostructures (active material, 70 wt%), polyvinylidene fluoride binder (10 wt%) and conductive carbon black (20 wt%), using *n*-methyl pyrrolidone as solvent to achieve a homogenous slurry, which was uniformly spread on Cu foil. The fabricated anodes were next dried at  $120^\circ\text{C}$  in order to remove traces of moisture and solvent, and then pressed to enhance the conductivity. Coin cells were prepared in an Ar atmosphere inside a glove box (M. Braun, USA). The electrodes were separated inside the coin cell using a Celgard 2400 membrane. The electrolyte consisted of lithium hexafluorophosphate ( $\text{LiPF}_6$ ) in ethylene carbonate and dimethyl carbonate (EC : DMC 1 : 1, v/v). All the electrochemical measurements were performed at room temperature using a Gamry Instruments potentiostat and PHE200 electrochemical software. The coin cells were tested in the potential window of 0.01–3.0 V vs.  $\text{Li/Li}^+$ , and the applied current density was  $200 \text{ mA cm}^{-2}$ . The ac impedance spectra were collected (after cell potential stabilization) by applying a sinusoidal signal with amplitude of 5.0 mV over the frequency range from 1 mHz to 100 kHz, and using EIS300 electrochemical software.

## Results and discussion

The XRD pattern of the as-synthesized MnS films is shown in Fig. 1a. The diffraction peaks observed were indexed to the diffraction planes of hexagonal wurtzite MnS ( $\gamma$ -MnS) phase [(100), (002), (101), (102), (110), (103), (200), (112) and (201)] according to the JCPDS card, File no. 40-1289 and space group

$P6_3mc$ . Based on the peak positions, the  $\alpha$ -MnS phase can be ruled out. However, the (002), (110) and (112) peaks observed overlap with (111), (220) and (311) peaks of  $\beta$ -MnS phase (JCPDS no. 401288), respectively, whose presence cannot be discarded. The simulation reveals that the main peak of the diffractogram consists of one less-intense peak (FWHM  $\sim 0.12$ ) attributable to (111) peak of  $\beta$ -MnS, and one more-pronounced sharp peak (FWHM  $\sim 0.24$ ) corresponding to (002) peak of  $\gamma$ -MnS (see inset of Fig. 1a). This simulation also shows that only an 8% of the area under (002) peak belongs to the  $\beta$ -MnS phase.

Similar findings were observed for (110) and (112) peaks of the diffractogram. The cubic  $\beta$ -MnS phase found is ascribed to the temperature gradient formed during the synthesis process, which is typical for MnS.<sup>6,24</sup> Otherwise, the sharp peaks found and the absence of impurities in the XRD pattern indicate that the metastable  $\gamma$ -MnS phase is of high crystalline quality and purity.

The  $\gamma$ -MnS films were also studied by Raman scattering spectroscopy (Fig. 1b). The well-pronounced low-frequency Raman band observed was simulated using the damped harmonic oscillator phonon model (DHOPM).<sup>6,9</sup> From the simulation, the Raman mode is determined to be at  $592\text{ cm}^{-1}$ , and can be ascribed to scattering involving two magnon branches (2M) occurring in MnS.<sup>25</sup> The corresponding fit shows that this band is red-shifted by  $\sim 23\text{ cm}^{-1}$  and narrow (FWHM  $\sim 22\text{ cm}^{-1}$ ) when compared to bulk MnS,<sup>25</sup> which is a clear indicator that the  $\gamma$ -MnS phase is under tensile stress developed during the growth process under CVD conditions. Note that this band is relatively weak owing to the fact that the Raman spectrum was collected at room temperature, which is above the Néel temperature ( $T_N \sim 90\text{ K}$ ) for  $\gamma$ -MnS nanostructures.<sup>26</sup> Taken altogether, the Raman analysis indicates the formation of MnS with wurtzite structure, consistent with the XRD analysis.

The FE-SEM images of  $\gamma$ -MnS taken at different magnifications are depicted in Fig. 2. The average thickness of  $\gamma$ -MnS films was found to be  $\sim 25\text{ }\mu\text{m}$  (see Fig. 2a and inset therein). These films present a densely populated surface with spaghetti-like NRs, with the catalytic material on the ends thereof (see Fig. 2b and inset therein), indicating that they grew following a catalyst-based mechanism. It was also found that  $\gamma$ -MnS NRs

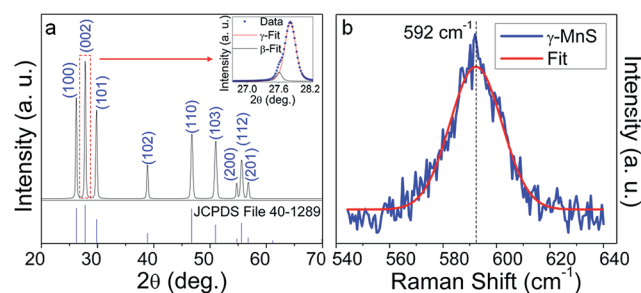


Fig. 1 (a) XRD pattern, and (b) micro-Raman spectrum of  $\gamma$ -MnS films. Inset in (a) shows the Voigt fit that evidences the co-existence of  $\beta$ -MnS in the products. Red solid line in (b) shows the DHOPM fit corresponding to the scattering Raman mode characteristic of  $\gamma$ -MnS.

possess a high aspect ratio (AR) with widths falling in the range of 100–350 nm and lengths over  $25\text{ }\mu\text{m}$ . Fig. 2c and d show that the majority of the NRs observed are comb-shaped, *i.e.*, NRs with one side flat and the other side saw-toothed. These NRs exhibiting this particular cross-section will be referred to from now on as NSs. Similar morphologies were recently reported by Ma *et al.* working on wurtzite ZnS and CdSe nanostructures.<sup>27,28</sup>  $\gamma$ -MnS NBs are also present in the as-grown samples, but they are not the dominant component. Extensive electron microscopy observations suggest that both saw- and belt-like nanostructures, (which were obtained in the same growth temperature range) possess a common nucleation site,<sup>27</sup> indicating that they are nucleated and grown as bunches randomly dispersed on the substrate.

The HRTEM image depicted in Fig. 3a shows that the geometrical shape of the  $\gamma$ -MnS nanostructures is belt-like with uniform width along their length and widths in the range of 80 to 250 nm. Catalytic particles were observed on the ends of the NBs, evidencing that the vapor–liquid–solid (VLS) process is the prevailing growth mechanism. The ripple-like contrast seen in inset of Fig. 3a is most likely a result of the strain generated from the bending of the belt, since the degree of flexibility of the NBs is associated to their thickness. However, it could also be related to the formation of stacking faults and twinning occurring as a mechanism to stabilize the growth of the nanostructures. It was also observed that various NBs grow from a single nucleation site, which is consistent with the SEM analysis. Statistical analyses of the HRTEM images reveal that the dominant morphology in the products is saw-like, such as that depicted in Fig. 3b. Individual  $\gamma$ -MnS NSs consist of one flat side and one side with extrusive teeth, which resemble parallel triangles ( $\sim 10$ – $50\text{ nm}$  high) grown along the NSs axes, similar to perfectly aligned nanocantilever arrays used in probing scanning technology. It was found that the ends (tips) of the cantilevers are atomically rough, as pointed out by the black arrowhead in the middle inset of Fig. 3b. The surfaces of the NSs

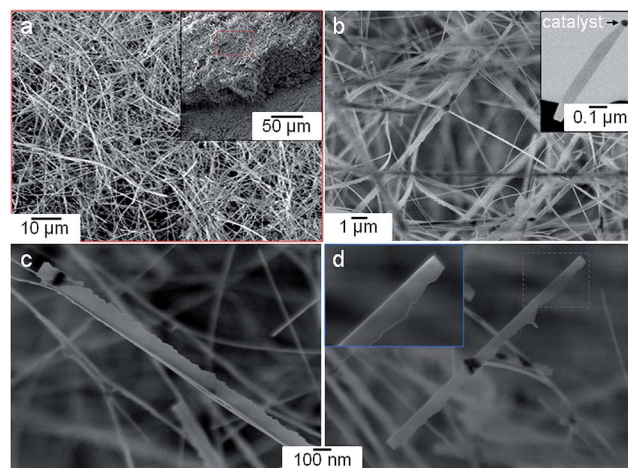


Fig. 2 (a–d) FE-SEM images of  $\gamma$ -MnS NSs. Inset in (a) shows a low-magnification FE-SEM image of the product. Insets in (b) and (d) show the catalytic and a close-up of an individual NS, respectively.

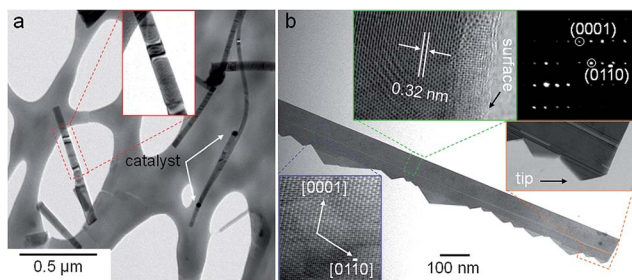


Fig. 3 (a) HRTEM image of  $\gamma$ -MnS NBs showing the catalyst employed in the synthesis. Inset: a closer view of an individual  $\gamma$ -MnS NB. (b) Representative HRTEM image of an individual  $\gamma$ -MnS NS. The left upper, middle, and lower insets in (b) show the surface, the tip and interior, respectively, of an individual  $\gamma$ -MnS NSs. The right upper inset in (b) shows the SAED pattern.

were clean, smooth, and atomically resolved (see left upper inset of Fig. 3b, that corresponds to the rotated green-delineated selected area in Fig. 3b). Hence, this method provides an efficient route to develop 1D building blocks without any sheathed amorphous phase, having low degree of interdiffusion between the material itself and the surrounding medium. An inner view into the nanostructures (see lower and left upper insets of Fig. 3b) and the selected area electron diffraction (SAED, operated in TEM mode) analysis (see right upper inset of Fig. 3b) corresponding to the selected area of left upper inset of Fig. 3b) evidence that: NSs grow along the  $[01\bar{1}0]$  with top and bottom surfaces  $\pm(2\bar{1}\bar{1}0)$ , and side surfaces  $(0001)$ ; the growth direction of the NSs teeth is along  $[0001]$  with an interplanar spacing of 0.32 nm; and they are single crystal with wurtzite structure.

In order to clearly establish the elemental composition of the products, we have conducted energy-dispersive X-ray spectroscopy (EDS) and electron energy loss spectroscopy (EELS) measurements on individual  $\gamma$ -MnS NSs. Both techniques were operated in STEM mode. The EDS line-scans depicted in Fig. 4b (carried out on the region shown in Fig. 4a) show that the nanostructures consist of Mn and S with a molar ratio of  $\sim 1 : 1$ . The width obtained by this technique ( $\sim 150$  nm) correlates well with the values range estimated from the HRTEM analysis. The EDS elemental mappings (not shown here) revealed the overall distribution of Mn and S in the nanostructures. Line-scanned EELS spectra are depicted in Fig. 4c, which were carried out on the region shown in Fig. 4a. The EELS spectra show a band peaking at  $\sim 55$  eV, which corresponds to the low-loss Mn  $M_{2,3}$  transition.<sup>6</sup> The spectra indicate that Mn fades at the teeth of the NSs, which is expected given that they are thinner than the backsaw (see next Section). The EELS probe used in this study was not adequate enough to resolve the light S atoms. However, a very weak signal was detected centered at  $\sim 165$  eV, which is attributed to the S  $L_{2,3}$  transition. Considered altogether, our findings show that single-crystal  $\gamma$ -MnS NSs were successfully synthesized by this method.

In order to gain more insight on the uniformity of the cross-sectional thickness, width of tip, and dislocations or defects present in  $\gamma$ -MnS NSs, we have carefully examined several NSs

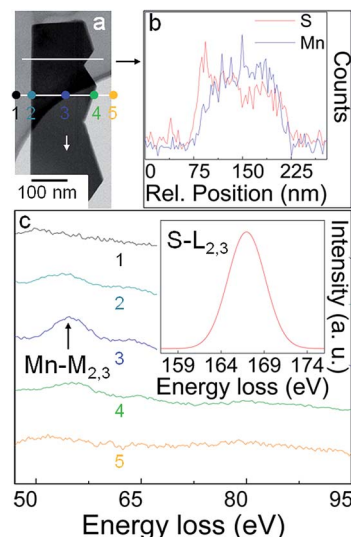


Fig. 4 (a) HRTEM image of an individual  $\gamma$ -MnS NS. (b and c) EDS profile and EELS spectra line-scanned from the region shown in (a). Inset in (c) shows the S  $L_{2,3}$ -shell edge of  $\gamma$ -MnS.

by illuminating them with 200 kV electrons for a somewhat longer time, but not enough to cause any increase of planar induced defects, damage, phase transformation or amorphous carbon deposition during the exposure. A representative HRTEM image of an individual NS is depicted in Fig. 5A. As noted, the NS is transparent to the TEM grid, which suggests that it is a few tens of nanometers thick. In fact, the backsaw thickness resulted to be in the range of 30 to 50 nm (see Fig. 5B and inset therein), and width-to-thickness ratios of  $\sim 3.3$  to 7.0. Interestingly, it was observed that the thickness was significantly reduced from the end of backsaw to the teeth (radially), reaching values as low as few nanometers. In this region, the NS (surface encompassed by the teeth) is structurally uniform and highly crystalline with two growth directions:  $[01\bar{1}0]$  and  $[0001]$ . In addition, no dislocations were found, but rather single atomic-thick stacking faults along the NS and teeth axes (see Fig. 5C and F). A closer view into the tip of teeth shows that its surface is seamless having a wurtzite structure (according to the fast Fourier transform (FFT) indexation), and its width is approximately a couple of atoms thick (see Fig. 5D and F). A unit-cell model for the hexagonal phase of MnS is shown in the inset of Fig. 5C. It was observed that the projected position of the Mn atoms in the unit cell correlates well with the TEM projection image taken on the end of tip and oriented along  $(2\bar{1}\bar{1}0)$ . The light sulfur atoms were less resolved by the TEM probe. Thus, the asymmetric growth behavior of the NR is ascribed to the polarization of the  $c$ -plane, which can be plausibly explained as follows: it is known that the chalcogen polar surfaces in chalcogenides and oxides (wurtzite ZnS, CdSe and ZnO) are unreactive,<sup>16,28,29</sup> and  $(0001)$ -Zn and  $-Cd$ -terminated surfaces are chemically reactive, specifically to adsorb methanol, formaldehyde and formic acid, and ZnS and CdSe vapors. It was found that these active surfaces are responsible for growing saw-like structures (such as one- and double-sized

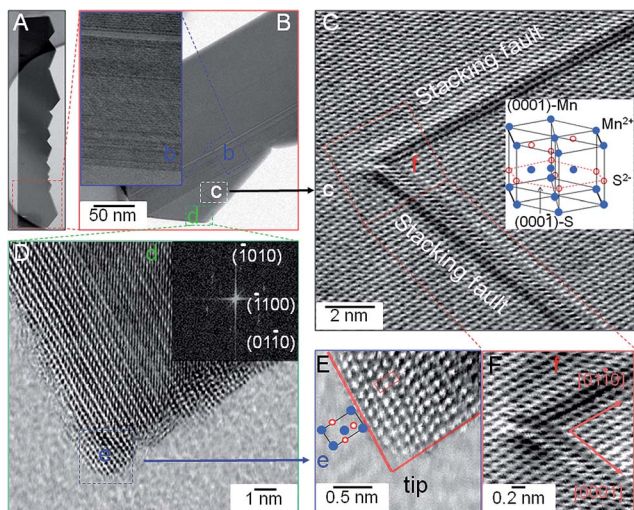


Fig. 5 (A) HRTEM image of a representative  $\gamma$ -MnS NS. (B) A closer view of (A) showing the interface between the backsaw and teeth (see inset therein). (C and F) Closer views into the teeth of (A) showing two stacking faults along NS and teeth axes. A unit-cell model of  $\gamma$ -MnS is shown in inset. (D and E) Lattice-resolved images of (A) showing the width of the nanotips. The FFT image corresponding to (D) is shown in inset of (D). Inset in (E) shows the projected position of the Mn atoms of the unit-cell model of  $\gamma$ -MnS shown in inset of (C).

NSs), acting as the self-catalyst.<sup>16,28,29</sup> Thus, polar surfaces induce directly the formation of asymmetric NRs. It is reasonable to claim that the saw-like nature of wurtzite-structured MnS consisting of alternating layers of oppositely charged ions ( $(0001)\text{-Mn}^{2+}$  and  $(000\bar{1})\text{-S}^{2-}$ ) and stacked parallel to the  $c$ -axis is induced by both the surface polarization and surface termination (see inset of Fig. 5C). Even though this ionic crystal composed of tetrahedrally-coordinated  $\text{Mn}^{2+}$  and  $\text{S}^{2-}$  layers exhibits an accumulating normal dipole moment (as a result of the surface polarization),  $\gamma$ -MnS NSs are shown to be stable without facets or massive surface reconstructions.

Taken altogether, the results indicate that single-crystal NSs were grown in two stages, as follows. First, the NRs are rapidly formed along  $[01\bar{1}0]$  with top and bottom surfaces  $\pm(2\bar{1}\bar{1}0)$ , and side surfaces  $(0001)$ , following a bottom-up Au-catalyst-based approach described by the VLS growth mechanism (see SEM analysis). Second, the subsequent growth (self-catalyzed from the Mn-terminated surface) of the teeth is conducted and oriented along  $[0001]$ , which is a consequence of the surface polarization and surface termination of the wurtzite structure (see HRTEM analysis).

We have evaluated the temperature-dependent PL properties of these NSs on Si substrate in order to provide further insight into their structural properties on a macroscopic scale, and evaluate their potential applicability as visible light emitters. The PL spectra of  $\gamma$ -MnS NSs recorded between  $T = 10$  to 300 K are depicted in Fig. 6a. An optical image of  $\gamma$ -MnS NSs suspended in aqueous solution as viewed under UV light excitation at 300 K is shown in the lower inset of Fig. 6a. The PL spectra at 300 K show one predominant yellowish-orange emission band peaking at  $\sim 578$  nm at, which is most likely ascribed to the

optically radiative recombination of electrons in shallow traps of surface localized states and photogenerated holes caused by stacking faults found in the NSs.<sup>30–32</sup> A relatively weak near band edge emission band centered at  $\sim 371$  nm and a defect-related emission band at 430 nm were observed at shorter wavelengths (see Fig. S1†). The near band edge emission of NSs is blueshifted  $\sim 16$  nm when compared to the bulk value (387 nm).<sup>19</sup> This shift is due possibly to the tiny width (a couple of atoms thick) of the  $\gamma$ -MnS NSs tips, which give rise to a quantum confinement effect. Similar results were reported for CdSe NSs.<sup>28</sup> However, another factor that may contribute to this blueshift is related to the small thickness of the NSs backsaws. The origin of the defect-related emission at 430 nm is presumably ascribed to the higher level transition of  $\gamma$ -MnS from the surface defects;<sup>22,33</sup> however, further research is needed to confirm such presumption. The PL intensity of the surface defect-induced emission at 582 nm is four orders of magnitude stronger than the exciton-related emission band at 371 nm. This confirms the presence of high density of single stacking faults in the NSs, which is consistent with the HRTEM analysis. As the temperature decreases (down to 60 K), the yellowish-orange emission band experiences a redshift of  $\sim 4$  nm, which is ascribed to its shallow-trap emission character (see upper inset of Fig. 6a). No substantial shift was detected below 60 K. Fig. 6b shows that the yellowish-orange emission band experiences a remarkable drop in PL intensity as the temperature decreases, as expected. A relatively small change in PL intensity was observed over the temperature range from 300 to 180 K. A significant increase in PL intensity was observed at lower temperatures. The integrated PL intensity was doubled from 100 to 10 K, which corresponds to the carrier transfer from electrons trapped in the donor states that are promoted towards the conduction band. It is assumed that the interaction between neighboring  $\text{Mn}^{2+}$  ions also contributes to the carrier transfer, which is influenced by the antiferromagnetic spin-ordering given that more prominent PL intensity changes were observed below 100 K ( $T_N \sim 90$  K).<sup>26,34</sup> To determine the position of the donor level with respect to the conduction band, the experimental data were simulated through a thermally activated carrier-transfer model<sup>35</sup> via

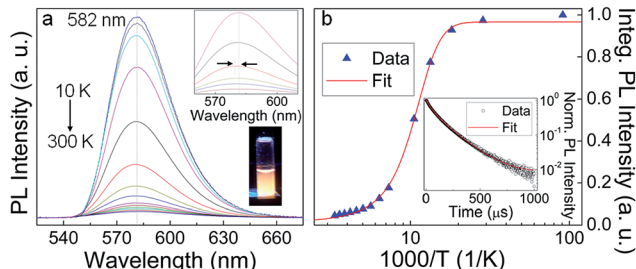


Fig. 6 (a) Temperature-dependent PL spectra of  $\gamma$ -MnS/Si films excited with a 325 nm He–Cd pulse laser. The redshift of the yellowish-orange emission band is shown in the upper inset. The lower inset shows the optical image of  $\gamma$ -MnS NSs dispersed in aqueous solution as viewed under UV light at 300 K. (b) Emission band intensity evolution as a function of temperature fitted to eqn (1). Inset in (b) shows the PL decay curve fitted with a bi-exponential function.

$$I = \frac{A}{1 + B \exp\left(-\frac{\varepsilon}{k_B T}\right)} \quad (1)$$

where  $A$ ,  $B$ ,  $\varepsilon$  and  $k_B$  are parameters related to the generation rate, the ratio of the nonradiative rate to the radiative rate, the activation energy for the thermal quenching, and the Boltzmann constant, respectively.<sup>35</sup> The value of the thermal activation energy from the fitting was  $\sim 40$  meV (comparable to that reported previously for MnS<sup>36</sup>), which signifies that the donor level involved is located  $\sim 40$  meV under the conduction band edge, which correlates well with the increased surface-to-volume ratio of  $\gamma$ -MnS NSs. This result supports the claim that the donor state is associated to the surface states of the NSs, and that the carrier transfer corresponds to a donor-acceptor pair transition.<sup>37,38</sup> To further understand the surface effect, we have examined the PL kinetics of the emission band at 582 nm. The PL decay curve at 300 K is depicted in the inset of Fig. 6b. Bi-exponential fitting shows that the long decay time component of these nanostructures was significantly reduced to  $\sim 44$   $\mu$ s, when compared to the bulk value of MnS, 140  $\mu$ s.<sup>34</sup> This evidences that the enhanced surface area of  $\gamma$ -MnS NSs is responsible for the acceleration of the non-radiative energy relaxation rate due to the energy transfer from the excited Mn-states to surface states. No significant changes in decay time were observed below  $T_N$ . Our large-scale observations indicate that these photoluminescent nanostructures remained structurally stable even at low temperatures.

As an example of the versatility of these nanostructures, we have evaluated their electrochemical properties as a potential and alternative anode material for LIBs. These experiments were motivated by the atomically defined morphology, highly stable structure and high degree of crystallinity that the nanostructures present, which play key roles in electrochemical energy storage. Other factors that influenced the use of MnS as an anode material were the great abundance and low cost of Mn and S, and the fact that MnS exhibits a theoretical specific capacity ( $616 \text{ mA h g}^{-1}$ ) higher than that of graphite ( $372 \text{ mA h g}^{-1}$ ), and possesses a high lithium storage capacity.<sup>6</sup> However, wurtzite-structured MnS as an anode material has been overlooked to date.

Charge-discharge experiments were carried out in coin cells, in the potential window of 0.01–3.0 V vs. Li/Li<sup>+</sup> and current density of  $200 \text{ mA g}^{-1}$ . The charge-discharge curves obtained are depicted in Fig. 7a. It was found that the initial specific discharge capacity and the corresponding charge capacity exceeded  $1150$  and  $710 \text{ mA h g}^{-1}$ , respectively. These values were found to be higher than those of bulk MnS, which is related to additional capacity resulting from reversible reactions that occur between displaced Mn and the electrolyte solvents.<sup>39</sup> The irreversible capacity loss (38%) is due presumably to the electrolyte decomposition, which would give rise to the subsequent formation of solid-electrolyte interphase (SEI) layers on the electrodes surface. Such SEI formation is common for most MnS anode materials.<sup>40,41</sup> A plateau was also observed at  $\sim 0.5$  V, which can be assigned to both carbon black additives and SEI layers. The extended flatness of this initial discharge voltage

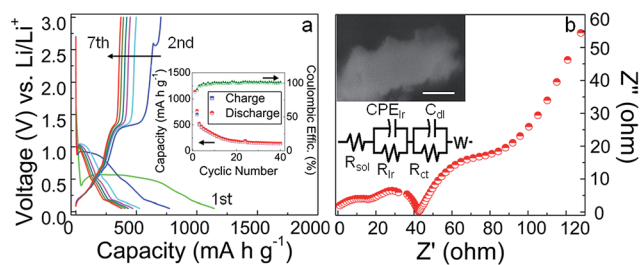
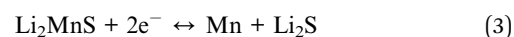
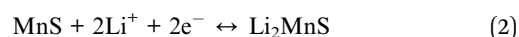


Fig. 7 (a) Charge-discharge voltage profiles and cyclic performance (inset), and (b) Nyquist plots of pristine  $\gamma$ -MnS NSs. Insets in (b) show the SEM image (scale bar: 100 nm) of an individual  $\gamma$ -MnS NSs after lithiation process and the equivalent circuit of the coin cell.

plateau is attributed to the superior crystallinity, stable surface site energetics, and larger electrochemical active surface area of  $\gamma$ -MnS NSs. From the second cycle onwards, this plateau vanishes almost completely signifying that the formation of the SEI layer takes place predominantly during the first cycle. The discharge capacity obtained in the second cycle was  $780 \text{ mA h g}^{-1}$ , representing a notable cyclic capacity retention  $\sim 70\%$ , being reduced up to 50% at the third cycle ( $530 \text{ mA h g}^{-1}$ ), from which point on small and regular changes were observed. This result is interesting given that the initial capacity of commercial MnS falls off severely over cycling.<sup>40</sup> In general, metal sulfides do not exhibit good charge cycling stability. Our results suggest that  $\gamma$ -MnS NWs show improved cycling with an increasing Coulombic efficiency of 96% in the second cycle, and retain a stable reversible capacity of  $\sim 200 \text{ mA h g}^{-1}$  after 40 cycles (see inset of Fig. 7a).

A rapid derivative of the capacity with respect to cell potential reveals three reduction peaks in the cathodic region: the peak at  $\sim 0.01$  V represents the lithiation of carbon;<sup>42</sup> the peak at  $\sim 0.2$  V is associated to the formation of Mn and Li<sub>2</sub>S (see eqn (3)), or to the non-aqueous electrolyte decomposition during the discharge process;<sup>40</sup> and the peak at  $\sim 0.7$  V is certainly attributed to the insertion of lithium into  $\gamma$ -MnS NSs structures to form a homogeneous phase of Li<sub>2</sub>MnS (see eqn (1)). Similarly, two oxidation zones are observed in the anodic region: the zone between 0.15–0.35 V corresponds to deintercalation of lithium stored in the carbon black;<sup>42</sup> and the zone between 0.6–1.3 V is ascribed to the extraction of lithium from  $\gamma$ -MnS.<sup>40,43</sup> These observations indicate that the SEI layer is thin and stable having high ionic conductivity, and the electrochemical reversible reactions in the MnS anode during the charge-discharge processes can be described by a two-stage lithiation mechanism as follows:<sup>40,44</sup>



To further investigate the interface reaction between the electrolyte solution and the anode electrode, we have conducted ac impedance measurements carried out at an open circuit potential of 0.6 V in the 23rd charge cycle. Nyquist plots depicted in Fig. 7b show two overlapping high-to-medium

frequency semicircles. The higher-frequency one is attributed both to the SEI layer at the active material/electrolyte interface and to the contact resistance, while the medium-frequency one corresponds to the charge transfer resistance on the electrode/electrolyte interface.<sup>39,41</sup> The spectrum in the low-frequency range is associated to synergistic resistive effects and to the Li-ions diffusion into the active material (Warburg resistance). The same results were obtained after 40 cycles suggesting that the anode material shows a good electrical contact between the active material and the current collector, and exhibits a good structural stability. These results can be summarized upon fitting the spectra of the cells and representing them as a first approximation by an equivalent circuit, as shown in inset of Fig. 7b, where  $R_{\text{sol}}$  is the electrolyte resistance, CPE is the constant phase element,  $R_{\text{lr}}$  is the surface layer-related resistance,  $C_{\text{dl}}$  is the double layer capacitance,  $R_{\text{ct}}$  is the charge transfer resistance at the interface and  $W$  is the Warburg element. It was also found that the diameters of the semicircles are relatively smaller than those of hollow cubic MnS microspheres and MnS nanoparticles.<sup>39,41</sup> This further evidences that  $\gamma$ -MnS electrode exhibits high electrical conductivity, which would induce in turn a rapid electron transport during the lithium insertion/extraction reaction and would enhance the reversible capacity.

Taken altogether, we suspect that the diffusion of  $\text{Li}^+$  ions towards the exposed teeth of MnS NSs provides abundant surface area for lithium diffusion through the interatomic spacing, and subsequent intercalations into different layers of the active material. As a result, a significant volumetric expansion (which is efficiently accommodated) and weak amorphization process (see inset of Fig. 7b) take place. Moreover, the structural flexibility of the NSs facilitates the prevention of detachment and agglomeration processes, thus allowing a more effective Li/electron diffusion during charge-discharge cycles. The improved cyclic capacity retention can be also assigned to the high degree of crystallinity, single-crystal architecture and structural integrity of these nanostructures, which offer transport advantages for both charge carriers.

## Conclusions

We have developed a synthetic method to produce high-quality dislocation-free single-crystal  $\gamma$ -MnS NSs by a straightforward CVD route, which employs lower synthesis temperatures (as compared to traditional thermal evaporation techniques  $\sim 1450$  K) and is scalable for mass production. These 1D building blocks possess a non-conventional cross-section, high aspect ratio, and a highly-stable surface structure. The crystals grow along the [0110] and [0001] directions and their surfaces are defined by unambiguous crystallographic planes. The asymmetric growth behavior is probably induced by the polarization of the  $c$ -plane, being the surface polarity a critical parameter to control the morphology of the nanostructures. This method employs  $\text{MnCl}_2$  and sulfur powder as starting materials instead of commercial  $\gamma$ -MnS powder, which would significantly reduce the production costs of  $\gamma$ -MnS and circumvent difficulties related to stoichiometry retention of  $\gamma$ -MnS powder on the

desired substrates during deposition process. These structurally uniform and geometrically perfect NSs are promising to be used as nanocantilever arrays for nanosensors and nanotweezers. This breakthrough advances novel methods to fabricate 1D building blocks suitable for optoelectronics and Li-ion battery applications.

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## Notes and references

- 1 Y. Xia, P. Yang, Y. Sun, Y. Wu, B. Mayers, B. Gates, Y. Yin, F. Kim and H. Yan, *Adv. Mater.*, 2003, **15**, 353–389.
- 2 Z. Deng, L. Tong, M. Flores, S. Lin, J. Cheng, H. Yan and Y. Liu, *J. Am. Chem. Soc.*, 2011, **133**, 5389–5396.
- 3 X. Duan, Y. Huang, Y. Cui, J. Wang and C. M. Lieber, *Nature*, 2001, **409**, 66–69.
- 4 Z. W. Pan, Z. R. Dai and Z. L. Wang, *Science*, 2001, **291**, 1947–1949.
- 5 N. G. Chopra, R. J. Luyken, K. Cherrey, V. H. Crespi, M. L. Cohen, S. G. Louie and A. Zettl, *Science*, 1995, **269**, 966–967.
- 6 J. Beltran-Huarac, O. Resto, J. Carpena-nuñez, W. M. Jadwisienzak, L. F. Fonseca, B. R. Weiner and G. Morell, *ACS Appl. Mater. Interfaces*, 2014, **6**, 1180–1186.
- 7 J. Ge, J. Wang, H. Zhang and Y. Li, *Chem.–Eur. J.*, 2004, **10**, 3525–3530.
- 8 F. Zhang and S. S. Wong, *Chem. Mater.*, 2009, **21**, 4541–4554.
- 9 J. Beltran-Huarac, J. Carpena-Nuñez, D. Barrionuevo, F. Mendoza, R. S. Katiyar, L. F. Fonseca, B. R. Weiner and G. Morell, *Carbon*, 2013, **65**, 252–260.
- 10 Z. L. Wang, *Annu. Rev. Phys. Chem.*, 2004, **55**, 159–196.
- 11 S. Liang, J. Zhou, G. Fang, J. Liu, Y. Tang, X. Li, A. Pan and L. Wang, *ACS Appl. Mater. Interfaces*, 2013, **5**, 8704–8709.
- 12 Z. L. Wang, *Adv. Mater.*, 2003, **15**, 432–436.
- 13 X. Fang, Y. Bando, G. Shen, C. Ye, U. K. Gautam, P. M. F. J. Costa, C. Zhi, C. Tang and D. Golberg, *Adv. Mater.*, 2007, **19**, 2593–2596.
- 14 X. Fang, Y. Bando, M. Liao, U. K. Gautam, C. Zhi, B. Dierre, B. Liu, T. Zhai, T. Sekiguchi, Y. Koide and D. Golberg, *Adv. Mater.*, 2009, **21**, 2034–2039.
- 15 X. J. Zheng, Y. Q. Chen, T. Zhang, B. Yang, C. B. Jiang, B. Yuan and Z. Zhu, *Sens. Actuators, B*, 2010, **147**, 442–446.
- 16 Z. Wang, L. L. Daemen, Y. Zhao, C. S. Zha, R. T. Downs, X. Wang, Z. L. Wang and R. J. Hemley, *Nat. Mater.*, 2005, **4**, 922–927.
- 17 S. Biswas, S. Kar and S. Chaudhuri, *Mater. Sci. Eng., B*, 2007, **142**, 69–77.
- 18 S. Wang, K. Li, R. Zhai, H. Wang, Y. Hou and H. Yan, *Mater. Chem. Phys.*, 2005, **91**, 298–300.



- 19 Y. Jun, Y. Jung and J. Cheon, *J. Am. Chem. Soc.*, 2002, **124**, 615–619.
- 20 J. Lu, P. Qi, Y. Peng, Z. Meng, Z. Yang, W. Yu and Y. Qian, *Chem. Mater.*, 2001, **13**, 2169–2172.
- 21 D. S. Kim, J. Y. Lee, C. W. Na, S. W. Yoon, S. Y. Kim and J. Park, *J. Phys. Chem. B*, 2006, **110**, 18262–18266.
- 22 Y. Ren, L. Gao, J. Sun, Y. Liu and X. Xie, *Ceram. Int.*, 2012, **38**, 875–881.
- 23 N. Moloto, N. J. Coville, S. S. Ray and M. J. Moloto, *Phys. B*, 2009, **404**, 4461–4465.
- 24 X. Yang, Y. Wang, K. Wang, Y. Sui, M. Zhang, B. Li, Y. Ma, B. Liu, G. Zou and B. Zou, *J. Phys. Chem. C*, 2012, **116**, 3292–3297.
- 25 H. Chou and H. Y. Fan, *Phys. Rev. B*, 1976, **13**, 3924–3938.
- 26 J. Banewicz and R. Lindsay, *Phys. Rev.*, 1956, **104**, 318–320.
- 27 C. Ma, D. Moore, J. Li and Z. L. Wang, *Adv. Mater.*, 2003, **15**, 228–231.
- 28 C. Ma, Y. Ding, D. Moore, X. Wang and Z. L. Wang, *J. Am. Chem. Soc.*, 2004, **126**, 708–709.
- 29 Z. L. Wang, X. Y. Kong and J. M. Zuo, *Phys. Rev. Lett.*, 2003, **91**, 185502.
- 30 H. Y. Wang, C. R. Wang, J. Xu, X. Liu, X. F. Xu, H. Z. Xing, L. J. Zhao and X. S. Chen, *J. Phys. D: Appl. Phys.*, 2012, **45**, 095301.
- 31 T. Tsuruoka, C. H. Liang, K. Terabe and T. Hasegawa, *Appl. Phys. Lett.*, 2008, **92**, 091908.
- 32 L. Yujie, Y. Liping, D. Ran, S. Pengbo, D. Honglin, Q. Yongping and Q. Guogang, *Nanotechnology*, 2004, **15**, 581–585.
- 33 J. Mu, Z. Gu, L. Wang, Z. Zhang, H. Sun and S. Kang, *J. Nanopart. Res.*, 2008, **10**, 197–201.
- 34 D. S. Kim, J. Y. Lee, C. W. Na, S. W. Yoon, S. Y. Kim, J. Park, Y. Jo and M. H. Jung, *J. Phys. Chem. B*, 2006, **110**, 18262–18266.
- 35 K. Kosai, B. T. Fitzpatrick, H. G. Grimmeiss, R. N. Bhargava and G. F. Neumak, *Appl. Phys. Lett.*, 1979, **35**, 194.
- 36 S. Varghese and M. Iype, *Orient. J. Chem.*, 2011, **27**, 265–269.
- 37 C. Ye, X. Fang, M. Wang and L. Zhang, *J. Appl. Phys.*, 2006, **99**, 063504.
- 38 J. Beltran-Huarac, J. Wang, H. Tanaka, W. M. Jadwisieniczak, B. R. Weiner and G. Morell, *J. Appl. Phys.*, 2013, **114**, 053106.
- 39 D. Chen, H. Quan, X. Luo and S. Luo, *Scr. Mater.*, 2014, **76**, 1–4.
- 40 Y. Liu, Y. Qiao, W. Zhang, Z. Li, X. Hu, L. Yuan and Y. Huang, *J. Mater. Chem.*, 2012, **22**, 24026–24033.
- 41 D. Chen, H. Quan, G. Wang and L. Guo, *ChemPlusChem*, 2013, **78**, 843–851.
- 42 B. J. Landi, M. J. Ganter, C. D. Cress, R. A. DiLeo and R. P. Raffaele, *Energy Environ. Sci.*, 2009, **2**, 638–654.
- 43 S. M. Lee, J. Lee and Y. C. Kang, *Chem.–Asian J.*, 2013, **9**, 590–595.
- 44 F. Liao, J. Światowska, V. Maurice, A. Seyeux, L. H. Klein, S. Zanna and P. Marcus, *Appl. Surf. Sci.*, 2013, **283**, 888–899.